

# High-Order Solutions and Generalized Darboux Transformations of Derivative Nonlinear Schrödinger Equations

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By using the bilinear method, we construct high-order solutions and generalized Darboux transformations of the derivative nonlinear Schrödinger equation (DNLS). The bilinear method is applied to the DNLS equation. As a consequence, we obtain the high-order solutions and the bilinear transformations of the DNLS equation.

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## 1. Introduction

The derivative nonlinear Schrödinger equation (DNLS) [1, 2]

$$u_t + u_{xx} + (|u|^2 u)_x = 0, \quad (1)$$

has a hierarchy of solutions, which can be obtained by the bilinear method. In this paper, we study the bilinear method for the DNLS equation. The bilinear method is applied to the DNLS equation. As a consequence, we obtain the high-order solutions and the bilinear transformations of the DNLS equation. The bilinear method is applied to the DNLS equation. As a consequence, we obtain the high-order solutions and the bilinear transformations of the DNLS equation.

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H gh-Vde VtV de c be the te ad V bā ee N VtV M e a  
 a t de b t ha g a at c a ch [4]. I the te Vg M e e  
 cate g t a fV d V (IST), the cV e V d tV t e- Ve VtV . I  
 the ca e M the K V t e eg-de V e e a V (KdV), the Ve t be e,  
 tha the ea V h gh-Vde V g a VtV dV V e t. I deed e  
 cV d M a t e- Ve VtV b Da bV t a fV d V (DT), ch  
 a VtV VtV [5]. The h gh-Vde VtV fV V ea Sch Vd ge  
 e a V (NLS) had bee t d ed b a a thV [4, 6, 7]. T V the be t M V  
 V edge, the h gh-Vde VtV M DNLS ha e e e bee e V t ed. The  
 a M th a e t V hV tha ch VtV a be M a ed b V-ca ed  
 ge e a ed Da bV t a fV d V (gDT).

Rece t , the Vg e a e he V e V [8], h ch a ea fV V he e  
 a d d a ea thV t at ace, ha bee a b ed M e te e t d . Th V e  
 a e , a V V a fea , V te , V ga t a e , a e cha a de ed th  
 age a t de a d M e a ea V the ea face. O e M the V be  
 a t V e a the Vg e a e the Vg e a e VtV a d V d a V  
 tab t a d the e a e e M V d V e b A h ed e ' g V [9 11].  
 D ffe e t a Vache ha e bee V V ed t V cV t d the ge e a ed Vg e  
 a e VtV M NLS, fV e a e , the age b V-ge V d c d hVd ad V t ed  
 b D ba d d a . [12, 13], O h a a d Ya g V the fa e V M H V a  
 b ea d hVd [14] h e the e e t a thV e the gDT a at V [15]. I  
 the IST te Vg , the h gh-Vde Vg e a e cV e V d tV t e- Ve  
 VtV a the b a ch Vt M the ed the V - a h g bac g V d  
 [16]. The t-Vde Vg e a e fV DNLS a M a ed b X a d cV V e  
 e e t [17]. H V e e , the h gh-Vde Vg e a e had e e bee t d ed.  
 We t ac eth V be b cV t d g gDT.

The DT [18, 19], h ch dV e V eed t V d V the e e ed a a a ,  
 V de a d ed a tV V e the La a e a V age b a ca . H V e e ,  
 the e a defed tha c a ca DT ca V be te a ed a the a e ed a  
 a a de . Th defed a e t V be t V cV t d the h gh-Vde Vg e  
 a e VtV b DT d ed . Th , e t V d f the DT d hVd. I th  
 V , e e te d the DT b the t tech e , M ha t a be te a ed a the  
 a e ed a a a de . The V d f ed t a fV d V e fe ed a gDT.

The e e cate g d hVd a ed t V t d DNLS th a h g  
 bac g V d (VBC) a d V - a h g bac g V d (NVBC) [3, 20 23]. The  
 N-b gh VtV fV a fV DNLS a e tab hed b Na a a a d Che  
 b the H V a b ea d hVd [24] a d the DT fV DNLS a cV t d ed b  
 I a [25] a d S e de [26] ( e e a V [17]). O e e a M th V the  
 cV t d V M the gDT a d ba ed V t , the h gh-Vde VtV VtV a d  
 Vg e a e a e M a ed. I add t V t V ab V e t V d M VtV , e  
 N- VtV a d h gh-Vde a V a VtV a e a V fV d.

based on the e-e-ta DT, e-cv-t-d fv KN the b-a DT, h ch  
 efe ed a the DT-II h e the e-e-ta DT efe ed a DT-I. We a V  
 te a e the e DT' a d V Vt the N-fvd DT' bWh fv DT-I a d DT-II,  
 a d cV de t V d ffe e t d M ed d V M the DT M the KN t V the  
 DNLS. I Sed V 3, the ge e a ed DT-I a d DT-II a e cV t d ed d a  
 b the t tech e. I Sed V 4, e cV de the a ca V M the gDT  
 a d ca c a e a V h gh-V de V t V, h ch c de h gh-V de b gh  
 V t V th the VBC, a d h gh-V de Vg e a e V t V. F a ed V  
 cV c de the a e a d Mfe V e d c V.

## 2. Darboux transformation for DNLS

Let us start with the following KN te [27]

$$\begin{aligned} u_t + u_{xx} + (u^2 v)_x &= 0, \\ -v_t + v_{xx} - (uv^2)_x &= 0, \end{aligned} \quad (2)$$

h ch a be te a the cV a b t cV dt V

$$U_t - V_x + [U, V] = 0, \quad (3)$$

M the ea te V La a

$$\Phi_x = U\Phi, \quad (4a)$$

$$\Phi_t = V\Phi, \quad (4b)$$

he e

$$U = -\frac{2}{\zeta^2}\sigma_3 + \frac{1}{\zeta}Q, \quad V = -\frac{2}{\zeta^4}\sigma_3 + \frac{2}{\zeta^3}Q - \frac{1}{\zeta^2}Q^2\sigma_3 + \frac{1}{\zeta}Q^3 - \frac{1}{\zeta}Q_x\sigma_3,$$

th

$$\sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}, \quad Q = \begin{pmatrix} 0 & u \\ -v & 0 \end{pmatrix}.$$

The e E a V (2) a e ed ced t M the DNLS (1) fv  $v = u^*$ . FV cV e e ce,  
 e t Vd ce the fv V g ad V t ea te fv (4),

$$-\Psi_x = \Psi U, \quad (5a)$$

$$-\Psi_t = \Psi V. \quad (5b)$$

### 2.1. DT-I

I the fv V g, e f t cV de the DT M the ed ced ea te (4).  
 Ge e a ea g, DT a ec a ga ge t a fv a V h ch ee the

for  $M$  is a constant. The explicit form of  $D$  is  $D = I + U$ , where  $U$  is a  $2 \times 2$  matrix depending on  $x$  and  $t$ , and  $I$  is the identity matrix.

$$D[1] = \zeta D_1 + D_0,$$

where  $D_1$  and  $D_0$  are  $2 \times 2$  matrices depending on  $x$  and  $t$ . The matrix  $D[1]$  is a DT matrix.

$$D[1]_x + D[1]U = U[1]D[1], \quad (d\mathbf{d}(D[1]))_x = T(U[1] - U)d\mathbf{d}(D[1]),$$

where  $U[1]$  is a  $2 \times 2$  matrix depending on  $x$  and  $t$ . After a gauge transformation, we find that the DT matrix  $D$  satisfies (4):

$$D[1] = \sigma_1 (\zeta + \zeta_1 - 2\zeta_1 P_1), \quad P_1 = \frac{\Phi_1 \Phi_1^T \sigma_1}{\Phi_1^T \sigma_1 \Phi_1}, \quad \sigma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}. \quad (6)$$

where  $\Phi_1 = (\varphi_1, \phi_1)^T$  is a vector depending on  $x$  and  $t$ . From (4) and  $\zeta = \zeta_1$ , and  $\Phi_1^T \sigma_1 \Phi_1$  is a scalar depending on  $x$  and  $t$ , we find that (5) and  $\zeta = -\zeta_1$ . Hence, we find that the DT matrix  $D$  satisfies the DNLS equation derived by Ablowitz [25]. With the help of the DT, Sode [26] and Xue et al. [17] calculated a vector-valued function for the DNLS. Next, we give a generalization of the above results for a DT.

**THEOREM 1.** *With  $\Phi_1$  and  $D[1]$  defined above,  $\Phi[1] = D[1]\Phi$  solves*

$$\Phi[1]_x = U[1]\Phi[1], \quad \Phi[1]_t =$$

From the identity (7). The identity for  $F_1(\zeta) \equiv D[1]_{,x} D[1]^{-1} + D[1] U D[1]^{-1} - U[1]$  at  $\zeta = \zeta_1$

$$\begin{aligned} \operatorname{Re}_{\zeta_1}(F_1(\zeta)) &= 2\zeta_1 \sigma_1 [-(I - P_1) P_{1,x} + (I - P_1) U(\zeta_1) P_1] \sigma_1 \\ &= -2\zeta_1 \sigma_1 (I - P_1) \Phi_1 \left[ \frac{\sigma_1 \Phi_1^T}{\sigma_1 \Phi_1^T \Phi_1} \right]_x \sigma_1 = 0, \end{aligned}$$

where we used the identity  $D[1]_{,x} D[1]^{-1} = -D[1](D[1]^{-1})_x$ . Similarly, the identity for  $F_1(\zeta)$  at  $\zeta = -\zeta_1$

$$\begin{aligned} \operatorname{Re}_{-\zeta_1}(F_1(\zeta)) &= -2\zeta_1 \sigma_1 [-P_{1,x}(I - P_1) + P_1 U(-\zeta_1)(I - P_1)] \sigma_1 \\ &= 2\zeta_1 \sigma_1 \left[ \frac{\Phi_1}{\sigma_1 \Phi_1^T \Phi_1} \right]_x \sigma_1 \Phi_1^T (I - P_1) \sigma_1 = 0. \end{aligned}$$

Define

$$U[1] = -\frac{1}{\zeta^2} \sigma_3 + \frac{1}{\zeta} Q[1],$$

and

$$Q[1] = \sigma_1 Q \sigma_1 - 2\zeta_1 \sigma_1 P_{1,x} \sigma_1, \quad (9)$$

the identity  $F_1(\zeta) D[1]$  is a rational function at  $\zeta = 0$ . The identity  $F_1(\zeta)$  is a rational function at  $\zeta = 0$ . It is easy to see that  $F_1(\zeta) \rightarrow 0$  as  $\zeta \rightarrow \infty$ . The identity (7) is a

Next, let us consider the identity (8). We will use the identity  $\widehat{V}[1] = -\frac{2}{\zeta^4} \sigma_3 + \frac{2}{\zeta^3} Q[1] + \frac{2}{\zeta^2} V_2 + \frac{1}{\zeta} V_1$ , that  $F_2(\zeta) \equiv D[1]_{,t} D[1]^{-1} + D[1] V D[1]^{-1} - \widehat{V}[1]$ . Proceeding as above, it follows that  $F_2(\zeta)$  is a rational function at  $\zeta = \pm\zeta_1$ , 0 and as  $\zeta \rightarrow \infty$ ,  $F_2(\zeta) \equiv 0$ .

In the following, we have  $\widehat{V}[1] = V[1]$ . Because of the identity at the identity  $(D[1]\Phi)_{xt} = (D[1]\Phi)_{tx}$ , we have

$$U[1]_t - \widehat{V}[1]_x + [U[1], \widehat{V}[1]] = 0. \quad (10)$$

Identifying the identity (10), we have

$$[\sigma_3, V_2] = 0, \quad (11)$$

$$[\sigma_3, V_1] = [V_2, Q[1]] + 2Q[1]_x, \quad (12)$$

$$[Q[1], V_1] = V_{2,x}. \quad (13)$$

From (11), we have  $V_2^{\text{eff}} = 0$ , where  $V_2^{\text{eff}}$  denotes the modified quantity at  $V_2$ . Similarly, we have

$$V_1^{\text{eff}} = \sigma_3 Q[1]_x - \frac{1}{2} \sigma_3 [Q[1], V_2] \quad (14)$$

through (12). Substituting (14) into (13) and using

$$V_2 = -Q[1]^2\sigma_3 + f(t).$$

Letting  $\zeta_1 = 0$ , we can read that  $V_2 = -\sigma_1 Q^2 \sigma_1 \sigma_3$ . The effective equation  $f(t) = 0$ . Moreover,  $V_1 = Q[1]^2 - Q[1]_x \sigma_3$ .

Since a gauge transformation  $D[1]^{-1}$  does not affect the admissible Lagrangian. Then the modified

*Remark 1.* In addition to (9) the effective equation for  $Q[1]$

$$Q[1] = \sigma_1(I - 2P_1)Q(I - 2P_1)\sigma_1 + \frac{2}{\zeta_1}\sigma_3\sigma_1(I - 2P_1)\sigma_1, \quad (15)$$

which are addressed in [17, 26]. We use (9) and (15), the effective equation

by the  $N$ -fold DT for the eeta DT (6), which effects a DT-I, the eeta

$$D[1] = \begin{pmatrix} -\zeta_1 \frac{\phi_+}{\phi_-} & \zeta \\ \zeta & \zeta_1 \frac{\phi_+}{\phi_-} \end{pmatrix}.$$

Adding the effective vector  $\Phi_i = (\varphi_i, \phi_i)^T$  of (4) and  $\zeta = \zeta_i (i = 1, 2, \dots, N)$  are given, the equation

PROPOSITION 1. [17, 25, 26] *The  $N$ -fold DT for DT-I can be represented as*

$$D_N = D[N]D[N-1]\cdots D[1] = \zeta^N \sigma_1^N + \sum_{k=0}^{N-1} \begin{pmatrix} \alpha_k & 0 \\ 0 & \beta_k \end{pmatrix} \sigma_1^k \zeta^k, \quad (16)$$

where  $\alpha_i$  are determined by the following equations

$$\begin{cases} \alpha_0 \varphi_i + \alpha_1 \zeta_i \phi_i + \cdots + \alpha_{2l-1} \zeta_i^{2l-1} \phi_i + \alpha_{2l} \zeta_i^{2l} \varphi_i = -\zeta_i^{2l+1} \phi_i, \\ \quad \text{he } N = 2l + 1; \\ \alpha_0 \varphi_i + \alpha_1 \zeta_i \phi_i + \cdots + \alpha_{2l-2} \zeta_i^{2l-2} \varphi_i + \alpha_{2l-1} \zeta_i^{2l-1} \phi_i = -\zeta_i^{2l} \varphi_i, \\ \quad \text{he } N = 2l, \end{cases}$$

$i = 1, 2, \dots, N$ . And  $\beta_i = \alpha_i(\varphi_j \leftrightarrow \phi_j)$ ,  $(i, j = 1, 2, \dots, N)$ .

The transformation between the fields is the following:

(i) When  $N = 2l + 1$

$$u[N] = -v - \left[ \frac{d\mathbf{d}(B)}{d\mathbf{d}(A)} \right]_x, \quad v[N] = -u + \left[ \frac{d\mathbf{d}(B(\varphi_j \leftrightarrow \phi_j))}{d\mathbf{d}(A(\varphi_j \leftrightarrow \phi_j))} \right]_x, \quad (17)$$

where  $A = (A_1^T, A_2^T, \dots, A_N^T)$ ,  $B = (B_1^T, B_2^T, \dots, B_N^T)$ ,

$$A_i = (\varphi_i, \zeta_i \phi_i, \dots, \zeta_i^{2l-2} \varphi_i, \zeta_i^{2l-1} \phi_i, \zeta_i^{2l} \varphi_i),$$

$$B_i = (\varphi_i, \zeta_i \phi_i, \dots, \zeta_i^{2l-2} \varphi_i, \zeta_i^{2l-1} \phi_i, \zeta_i^{2l+1} \phi_i).$$

(ii) When  $N = 2l$

$$u[N] = u - \left[ \frac{d\mathfrak{d}(D)}{d\mathfrak{d}(C)} \right]_x, \quad v[N] = v + \left[ \frac{d\mathfrak{d}(D(\varphi_j \leftrightarrow \phi_j))}{d\mathfrak{d}(C(\varphi_j \leftrightarrow \phi_j))} \right]_x, \quad (18)$$

where  $C = (C_1^T, C_2^T, \dots, C_N^T)$ ,  $D = (D_1^T, D_2^T, \dots, D_N^T)$ ,

$$C_i = (\varphi_i, \zeta_i \phi_i, \dots, \zeta_i^{2l-3} \phi_i, \zeta_i^{2l-2} \varphi_i, \zeta_i^{2l-1} \phi_i),$$

$$D_i = (\varphi_i, \zeta_i \phi_i, \dots, \zeta_i^{2l-3} \phi_i, \zeta_i^{2l-2} \varphi_i, \zeta_i^{2l} \phi_i).$$

## 2.2. DT-II

In this section, we show that the two-soliton solution of the DT-II equation [28, 29], derived by DT-II, is a particular case of the DT-I equation. For convenience, we let  $D[1]$  and  $D[1]^{-1}$  and  $\Psi_1$  be

$$D[1] = \begin{pmatrix} -\zeta_1 \frac{\varphi_1}{\phi_1} & \zeta \\ \zeta & -\zeta_1 \frac{\phi_1}{\varphi_1} \end{pmatrix}, \quad D[1]^{-1} = \frac{1}{\zeta^2 - \zeta_1^2} \begin{pmatrix} \zeta_1 \frac{\phi_1}{\varphi_1} & \zeta \\ \zeta & \zeta_1 \frac{\varphi_1}{\phi_1} \end{pmatrix}. \quad (19)$$

Since we have the vector  $\Psi_1 = (\chi_1, \psi_1)$  from the adjoint equation (5) at  $\zeta = \xi_1$  and the vector  $\Psi_1[1] = \Psi_1 D[1]^{-1}|_{\zeta=\xi_1}$  and the vector  $\Psi_1[1]$  from the adjoint equation (5) at  $\zeta = \xi_1$ . It can be seen that  $\sigma_1 \Psi_1[1]^T$  is a vector from the adjoint equation (5) at  $\zeta = -\xi_1$ . Therefore, we can deduce that the adjoint equation (5) at  $\zeta = -\xi_1$  is  $\sigma_1 \Psi_1[1]^T$ . Based on this, we get the adjoint equation (5) at  $\zeta = -\xi_1$  is  $\sigma_1 \Psi_1[1]^T$ .

$$T[1] = I + \frac{A}{\zeta - \xi_1} - \frac{\sigma_3 A \sigma_3}{\zeta + \xi_1}, \quad A = \frac{\xi_1^2 - \zeta_1^2}{2} \begin{pmatrix} \alpha & 0 \\ 0 & \beta \end{pmatrix} \Phi_1 \Psi_1, \quad (20)$$

where

$$\alpha^{-1} = \Psi_1 \begin{pmatrix} \xi_1 & 0 \\ 0 & \zeta_1 \end{pmatrix} \Phi_1, \quad \beta^{-1} = \Psi_1 \begin{pmatrix} \zeta_1 & 0 \\ 0 & \xi_1 \end{pmatrix} \Phi_1.$$

Furthermore, we have

$$T[1]^{-1} = I + \frac{B}{\zeta - \xi_1} - \frac{\sigma_3 B \sigma_3}{\zeta + \xi_1}, \quad B = \frac{\zeta_1^2 - \xi_1^2}{2} \Phi_1 \Psi_1 \begin{pmatrix} \beta & 0 \\ 0 & \alpha \end{pmatrix}. \quad (21)$$

Then a  $\mathbb{C}[v]$ -algebra  $\widehat{Q}$  and a  $\mathbb{C}[v]$ -algebra  $\widehat{Q}$

$$\widehat{Q} = Q + (A - \sigma_3 A \sigma_3)_x. \quad (22)$$

AbVe d c. V d ca e tha  $T[1]$ , a e the DT-II, de ed at V f V d DT f V  $D[1]$  the ca e M DNLS. We e a tha f V the t V c V V e t DNLS the a a V g M DT-II e t [30] h e the c V e V d g DT-I ha V bee c V t d ed.

I ha f V V, e c V de the te a V M the DT-II. A e tha e ha e  $N$  d t d V t V  $\Phi_i(\mu_i) = (\phi_i, \phi_i)$  M (4) a  $\zeta = \mu_i$  a d  $N$  d t d V t V  $\Psi_i(v_i) = (\chi_i, \psi_i)$  M (5) a  $\zeta = v_i$ . S a t V DT-I, e V th DT-II a d ha e the f V V g V V t V

PROPOSITION 2. *The  $N$ -fold DT for the DT-II could be written as the following form*

$$T_N = T[N]T[N-1] \cdots T[1] = I + \sum_{i=1}^N \left( \frac{C_i}{\zeta - v_i} - \frac{\sigma_3 C_i \sigma_3}{\zeta + v_i} \right), \quad (23)$$

and

$$T_N^{-1} = T[1]^{-1}T[2]^{-1} \cdots T[N]^{-1} = I + \sum_{i=1}^N \left( \frac{D_i}{\zeta - \mu_i} - \frac{\sigma_3 D_i \sigma_3}{\zeta + \mu_i} \right). \quad (24)$$

*Proof:* We ca c a the e d e f V b V h de V M (23)

Re  $|_{\zeta=v_i}(T_N)$

$$= \left( I + \frac{A_N}{v_i - v_N} - \frac{\sigma_3 A_N \sigma_3}{v_i + v_N} \right) \cdots A_i \cdots \left( I + \frac{A_1}{v_i - v_1} - \frac{\sigma_3 A_1 \sigma_3}{v_i + v_1} \right),$$

Re  $|_{\zeta=-v_i}(T_N)$

$$= - \left( I + \frac{A_N}{-v_i - v_N} - \frac{\sigma_3 A_N \sigma_3}{-v_i + v_N} \right) \cdots \sigma_3 A_i \sigma_3 \cdots \left( I + \frac{A_1}{-v_i - v_1} - \frac{\sigma_3 A_1 \sigma_3}{-v_i + v_1} \right).$$

Beca e V Re  $|_{\zeta=v_i}(T_N) = -\sigma_3 \text{Re } |_{\zeta=-v_i}(T_N) \sigma_3$ , E a V (23) a d. S a, (24) ca be V ed. ■

The  $N$ -f V d DT-II  $T_N$  a V t V r d the t a f V a V b e e the e d  $u, v$  a d  $u[N], v[N]$ , h ch a e g e be V

THEOREM 2. *The  $N$ -fold DT-II  $T_N$  induces the following transformations for the fields*

$$u[N] = u - 2 \left( \frac{d \mathbf{d} M_1}{d \mathbf{d} M} \right)_x, \quad v[N] = v + 2 \left( \frac{d \mathbf{d} N_1}{d \mathbf{d} N} \right)_x, \quad (25)$$

where  $M = (M_{ij})_{N \times N}$ ,  $N = (N_{ij})_{N \times N}$ ,  $M_{ij} = \frac{\Psi_i \sigma_3 \Phi_j}{\mu_j + \nu_i} - \frac{\Psi_i \Phi_j}{\mu_j - \nu_i}$ ,  $N_{ij} = -[\frac{\Psi_i \sigma_3 \Phi_j}{\mu_j + \nu_i} + \frac{\Psi_i \Phi_j}{\mu_j - \nu_i}]$ ,

$$M_1 = \begin{pmatrix} M_{11} & M_{12} & \cdots & M_{1N} & \psi_1 \\ M_{21} & M_{22} & \cdots & M_{2N} & \psi_2 \\ \vdots & \vdots & \ddots & \vdots & \vdots \\ M_{N1} & M_{N2} & \cdots & M_{NN} & \psi_{N_i} \\ \phi_1 & \phi_2 & \cdots & \phi_N & 0 \end{pmatrix}, \quad N_1 = \begin{pmatrix} N_{11} & N_{12} & \cdots & N_{1N} & \chi_1 \\ N_{21} & N_{22} & \cdots & N_{2N} & \chi_2 \\ \vdots & \vdots & \ddots & \vdots & \infty \\ & & & \mathbf{N} & \end{pmatrix}$$



Let  $\Psi$  be a solution of DT-I and DT-II associated with  $\Psi$  of DNLS. Hence, the DT-I of (28) can be extended to the  $N$ -band by the  $N$ -band  $\Psi$  of DNLS. The DT-II of (29) can be extended to the  $N$ -band by the  $N$ -band  $\Psi$  of DNLS.

### 3. Generalized Darboux transformations

In this section, we extend the  $gDT$  associated with  $D[1]$  and  $T[1]$ . We first consider the  $gDT$  associated with the  $N$ -band  $\Psi$  [15]. Indeed, the  $gDT$  associated with  $\Psi$  can be extended to the  $N$ -band  $\Psi$  of DNLS. The  $gDT$  associated with  $\Psi$  can be extended to the  $N$ -band  $\Psi$  of DNLS.

$$\Phi_1^{[1]} = \lim_{\epsilon \rightarrow 0} \frac{(D[1]\Phi_1)|_{\zeta=\zeta_1+\epsilon}}{\epsilon},$$

where

$$\Phi_1^{[1]} = \lim_{\epsilon \rightarrow 0} \frac{(T[1]\Phi_1)|_{\zeta=\zeta_1+\epsilon}}{\epsilon},$$

which is the extended  $gDT$  associated with  $\Psi$ .

#### 3.1. $gDT$ -I

Let  $\Psi$  be a solution of the  $gDT$  associated with DT-I, then we have  $n$   $N$ -band  $(\varphi_i, \phi_i)^T$  associated with the  $N$ -band  $\Psi$  of DNLS. For  $i = 1, \dots, n$ , the  $gDT$  associated with  $\Psi$  can be extended to the  $N$ -band  $\Psi$  of DNLS.

$$D_1^{[0]} = \begin{pmatrix} -\zeta_1 \frac{\varphi_1}{\phi_1} & \zeta \\ \zeta & -\zeta_1 \frac{\phi_1}{\varphi_1} \end{pmatrix}.$$

Associated with  $\Psi$ , the  $gDT$  associated with  $\Psi$  can be extended to the  $N$ -band  $\Psi$  of DNLS.

$$\begin{aligned} \begin{pmatrix} \varphi_1^{[1]} \\ \phi_1^{[1]} \end{pmatrix} &= \lim_{\epsilon \rightarrow 0} \frac{D_1^{[0]}|_{\zeta=\zeta_1} + \epsilon \sigma_1}{\epsilon} \begin{pmatrix} \varphi_1(\zeta_1 + \epsilon) \\ \phi_1(\zeta_1 + \epsilon) \end{pmatrix} \\ &= D_1^{[0]}|_{\zeta=\zeta_1} \frac{d}{d\zeta} \begin{pmatrix} \varphi_1(\zeta) \\ \phi_1(\zeta) \end{pmatrix} \Big|_{\zeta=\zeta_1} + \sigma_1 \begin{pmatrix} \varphi_1(\zeta_1) \\ \phi_1(\zeta_1) \end{pmatrix} \end{aligned}$$

where  $u = u[1]$  and  $v = v[1]$  at  $\zeta = \zeta_1$ , which is the extended  $gDT$  associated with  $\Psi$ .

$$D_1^{[1]} = \begin{pmatrix} -\zeta_1 \frac{\varphi_1^{[1]}}{\phi_1^{[1]}} & \zeta \\ \zeta & -\zeta_1 \frac{\phi_1^{[1]}}{\varphi_1^{[1]}} \end{pmatrix}.$$

Then we can verify that the following holds

**THEOREM 3.** *Let  $(\varphi_i, \phi_i)^T$  be the solutions of Lax pair at  $\zeta = \zeta_i$  ( $i = 1, \dots, n$ ) and assume that DT-I possesses  $m_i$  order zeros at  $\zeta = \zeta_i$ . Then we have the following gDT-I:*

$$D_N = D_n^{[m_n-1]} \dots D_n^{[1]} D_n^{[0]} \dots D_1^{[m_1-1]} \dots D_1^{[1]} D_1^{[0]}, \quad (30)$$

where

$$N = \sum_{i=1}^n m_i,$$

and

$$D_i^{[j]} = \begin{pmatrix} -\zeta_i \frac{\varphi_i^{[j-1]}}{\phi_i^{[j-1]}} & \zeta \\ \zeta & -\zeta_i \frac{\phi_i^{[j-1]}}{\varphi_i^{[j-1]}} \end{pmatrix},$$

$$\begin{pmatrix} \varphi_i^{[j-1]} \\ \phi_i^{[j-1]} \end{pmatrix} = \sum_{l=1}^{j-1} \frac{\Omega_{j-1-l}}{l!} \frac{d}{d\zeta^l} \begin{pmatrix} \varphi_i \\ \phi_i \end{pmatrix} \Big|_{\zeta=\zeta_i}, \quad \begin{pmatrix} \varphi_i^{[0]} \\ \phi_i^{[0]} \end{pmatrix} = \begin{pmatrix} \varphi_i \\ \phi_i \end{pmatrix}, \quad (31)$$

and

$$\Omega_l = \sum_{\sum \delta_i^k = l} M_i^{[j-2]} \dots M_i^{[0]} \dots M_1^{[m_1-1]} \dots M_1^{[0]}, \quad M_i^{[k]} = \begin{cases} \sigma_1, & \text{if } \delta_i^k = 1; \\ D_i^{[k]}|_{\zeta=\zeta_i}, & \text{if } \delta_i^k = 0. \end{cases}$$

*Proof:* To verify that the gDT-I, we take the following DT

$$D_1^{[0]} = \begin{pmatrix} -\zeta_1 \frac{\varphi_1}{\phi_1} & \zeta \\ \zeta & -\zeta_1 \frac{\phi_1}{\varphi_1} \end{pmatrix}.$$

By the way, we take a vector  $(\varphi_1[1], \phi_1[1])$ , and the vector  $(\varphi_i, \phi_i)^T$  is a factor of  $D_1^{[1]}$ . To get the general form of the vector  $(\varphi_i, \phi_i)^T$ , we find the following

$$\begin{pmatrix} \varphi_i^{[j]} \\ \phi_i^{[j]} \end{pmatrix} = \lim_{\epsilon \rightarrow 0} \frac{[D_i^{[j-1]} \dots D_i^{[1]} D_i^{[0]} \dots D_1^{[m_1-1]} \dots D_1^{[1]} D_1^{[0]}] \Big|_{\zeta=\zeta_i+\epsilon}}{\epsilon^j} \begin{pmatrix} \varphi_i(\zeta_i + \epsilon) \\ \phi_i(\zeta_i + \epsilon) \end{pmatrix},$$

which is the vector we are interested in. The case of the vector  $(\varphi_i, \phi_i)^T$  is similar.  $\blacksquare$

To have a vector added at the end of the gDT-I, we take the vector  $(\varphi_i, \phi_i)^T$  and the  $N$ -fold DT-I (16). It follows from (17) and (18) that the vector  $(\varphi_i, \phi_i)^T$  is the red one:

( ) Whe  $N = 2l + 1$

$$u[N] = -v - \left[ \frac{d\mathfrak{A}(B)}{d\mathfrak{A}(A)} \right]_x, \quad v[N] = -u + \left[ \frac{d\mathfrak{A}(B(\varphi_j \leftrightarrow \phi_j))}{d\mathfrak{A}(A(\varphi_j \leftrightarrow \phi_j))} \right]_x, \quad (32)$$

he e

$$A =$$

$\mu = \mu_i$  and  $\Psi_i(\zeta = v) = (\chi_i, \psi_i)$  are given by the adjoint Lax pair (5) at  $v = v_i$  ( $i = 1, \dots, n$ ).

**THEOREM 4.** Let  $\Phi_i$  be the solutions of Lax pair (4) at  $\zeta = \mu_i$  and  $\Psi_i$  be the solutions of adjoint Lax pair (5) at  $\zeta = v_i$  ( $i = 1, \dots, n$ ),

$$\sum_{i=1}^r m_i = N,$$

assume that DT-II possesses  $m_i$  order zeros at  $\zeta = \pm\mu_i$  and inverse of DT-II possesses  $m_i$  order zeros at  $\zeta = \pm v_i$ . Then we have the following gDT-II

$$\begin{aligned} T_N &= T_n^{[m_i-1]} \dots T_n^{[0]} \dots T_1^{[m_1-1]} \dots T_1^{[0]}, \\ T_N^{-1} &= (T_1^{[0]})^{-1} \dots (T_1^{[m_1-1]})^{-1} \dots (T_n^{[0]})^{-1} \dots (T_n^{[m_i-1]})^{-1} \end{aligned} \quad (34)$$

where

$$\begin{aligned} T_i^{[j]} &= I + \frac{A_i^{[j]}}{\zeta - v_i} - \frac{\sigma_3 A_i^{[j]} \sigma_3}{\zeta + v_i}, \quad (T_i^{[j]})^{-1} = I + \frac{B_i^{[j]}}{\zeta - \mu_i} - \frac{\sigma_3 B_i^{[j]} \sigma_3}{\zeta + \mu_i}, \\ A_i^{[j]} &= \frac{v_i^2 - \mu_i^2}{2} \begin{pmatrix} \alpha_i^{[j]} & 0 \\ 0 & \beta_i^{[j]} \end{pmatrix} \Phi_i^{[j]} \Psi_i^{[j]}, \quad B_i^{[j]} = \frac{\mu_i^2 - v_i^2}{2} \Phi_i^{[j]} \Psi_i^{[j]} \begin{pmatrix} \beta_i^{[j]} & 0 \\ 0 & \alpha_i^{[j]} \end{pmatrix}, \\ (\alpha_i^{[j]})^{-1} &= \Psi_i^{[j]} \begin{pmatrix} v_i & 0 \\ 0 & \mu_i \end{pmatrix} \Phi_i^{[j]}, \quad (\beta_i^{[j]})^{-1} = \Psi_i^{[j]} \begin{pmatrix} \mu_i & 0 \\ 0 & v_i \end{pmatrix} \Phi_i^{[j]}, \end{aligned}$$

and

$$\Phi_i^{[j]} = \sum_{l=0}^j \frac{\Omega_l}{(j-l)!} \frac{d^{j-l}}{d\mu^{j-l}} \Phi_i|_{\mu=\mu_i}, \quad \Omega_l = \sum_{\sum \delta_i^j = l} M_i^{[j-1]} \dots M_i^{[0]} \dots M_1^{[m_1-1]} \dots M_1^{[0]},$$

$$\Psi_i^{[j]} = \sum_{l=0}^j \frac{d^{j-l}}{dv^{j-l}} \Psi_i|_{v=v_i} \frac{\Lambda_l}{(j-l)!}, \quad \Lambda_l = \sum_{\sum \delta_i^j = l} N_i^{[j-1]} \dots N_i^{[0]} \dots N_1^{[m_1-1]} \dots N_1^{[0]},$$

$$M_i^{[j]} = \begin{cases} \frac{1}{\mu_i^2 - v_i^2}, & \text{if } \delta_i^j = 2; \\ \frac{2\mu_i + A_i^{[j]} - \sigma_3 A_i^{[j]} \sigma_3}{\mu_i^2 - v_i^2}, & \text{if } \delta_i^j = 1; \\ T_i^{[j]}|_{\zeta=\mu_i}, & \text{if } \delta_i^j = 0. \end{cases}, \quad N_i^{[j]} = \begin{cases} \frac{1}{v_i^2 - \mu_i^2}, & \text{if } \delta_i^j = 2; \\ \frac{2v_i + B_i^{[j]} - \sigma_3 B_i^{[j]} \sigma_3}{v_i^2 - \mu_i^2}, & \text{if } \delta_i^j = 1; \\ (T_i^{[j]})^{-1}|_{\zeta=v_i}, & \text{if } \delta_i^j = 0. \end{cases}$$

*Proof:* Next, let us show that the DT-II gauge (20) and the Lax pair (4)  $(\Phi[1], U[1], V[1])$  at  $\zeta = \mu_1$  and at  $\zeta = v_1$  are (5)  $(\Psi[1], U[1], V[1])$  at  $\zeta = v_1$ , that is,

$$\begin{aligned}\Phi_1^{[1]} &= \lim_{\delta \rightarrow 0} \frac{T_1^{[0]}|_{\zeta=\mu_1+\delta} \Phi_1(\mu_1 + \delta)}{\delta} = T_1^{[0]}|_{\zeta=\mu_1} \frac{d}{d\mu} \Phi_1|_{\mu=\mu_1} + S_1 \Phi_1(\mu_1), \\ \Psi_1^{[1]} &= \lim_{\delta \rightarrow 0} \frac{\Psi_1(v_1 + \delta) T_1^{[0]-1}|_{\zeta=v_1+\delta}}{\delta} = \frac{d}{dv} \Psi_1|_{v=v_1} T_1^{[0]-1}|_{\zeta=v_1} + \Psi_1(v_1) R_1,\end{aligned}$$

where

$$S_1 = \frac{2\mu_1 + A_1 - \sigma_3 A_1 \sigma_3}{\mu_1^2 - v_1^2}, \quad R_1 = \frac{2v_1 + B_1 - \sigma_3 B_1 \sigma_3}{v_1^2 - \mu_1^2}.$$

Therefore, we can write the Lax pair of the DT-II  $T[2]$

$$T_1^{[1]} = I + \frac{A_1^{[1]}}{\zeta - v_1} - \frac{\sigma_3 A_1^{[1]} \sigma_3}{\zeta + v_1}, \quad (T_1^{[1]})^{-1} = I + \frac{B_1^{[1]}}{\zeta - \mu_1} - \frac{\sigma_3 B_1^{[1]} \sigma_3}{\zeta + \mu_1}.$$

Given a target accuracy, if the gauge  $\Phi_i$  and  $\Psi_i$ , we can find the

$$\begin{aligned}\Phi_i^{[j]} &= \lim_{\delta \rightarrow 0} \frac{(T_i^{[j-1]} \dots T_i^{[0]} \dots T_1^{[m_1-1]} \dots T_1^{[0]})|_{\zeta=\mu_i+\delta}}{\delta^j} \Phi_i(\mu_i + \delta), \\ \Psi_i^{[j]} &= \lim_{\delta \rightarrow 0} \Psi_i(v_i + \delta) \frac{(T_1^{[0]-1} \dots (T_1^{[m_1-1]})^{-1} \dots (T_i^{[0]})^{-1} \dots (T_i^{[j-1]})^{-1})|_{\zeta=v_i+\delta}}{\delta^j},\end{aligned}$$

and then calculate the DT-II (34). The case of the

Defining the Lax pair, then a few details are

$$u[N] = u + 2 \sum_{i=1}^n \sum_{j=0}^{m_i-1} [(A_i^{[j]})_{12}]_x, \quad (35)$$

$$v[N] = v - 2 \sum_{i=1}^n \sum_{j=0}^{m_i-1} [(A_i^{[j]})_{21}]_x. \quad (36)$$

As before, the Lax pair (35) and (36) can be written as

$$u[N] = u - 2 \left( \frac{d(P_1)}{d(P)} \right)_x, \quad v[N] = v + 2 \left( \frac{d(Q_1)}{d(Q)} \right)_x \quad (37)$$

he e

$$P_1 = \begin{pmatrix} P^{[11]} & P^{[12]} & \dots & P^{[1r]} & \widehat{\psi}_1 \\ P^{[21]} & P^{[22]} & \dots & P^{[2r]} & \widehat{\psi}_2 \\ \vdots & \vdots & \ddots & \vdots & \vdots \\ P^{[r1]} & P^{[r2]} & \dots & P^{[rr]} & \widehat{\psi}_r \\ \widehat{\varphi}_1 & \widehat{\varphi}_2 & \dots & \widehat{\varphi}_r & 0 \end{pmatrix}, \quad P = \begin{pmatrix} P^{[11]} & P^{[12]} & \dots & P^{[1r]} \\ P^{[21]} & P^{[22]} & \dots & P^{[2r]} \\ \vdots & \vdots & \ddots & \vdots \\ P^{[r1]} & P^{[r2]} & \dots & P^{[rr]} \end{pmatrix},$$

th

$$\widehat{\psi}_i = \left( \psi_i, \frac{\partial}{\partial v} \psi_i, \dots, \frac{1}{(m_i - 1)!} \frac{\partial^{m_i-1}}{\partial v^{m_i-1}} \psi_i \right)^T \Big|_{v=v_j},$$

$$\widehat{\varphi}_j = \left( \varphi_j, \frac{\partial}{\partial \mu} \varphi_j, \dots, \frac{1}{(m_j - 1)!} \frac{\partial^{m_j-1}}{\partial \mu^{m_j-1}} \varphi_j \right) \Big|_{\mu=\mu_i},$$

$$P^{[ij]} = \left( P_{kl}^{[ij]} \right)_{m_k, m_l},$$

$$P_{kl}^{[ij]} = \frac{1}{(k-1)!(l-1)!} \frac{\partial^{k+l-2}}{\partial v^{k-1} \partial \mu^{l-1}} \left( \frac{\Psi_i(v) \sigma_3 \Phi_j(\mu)}{\mu + v} - \frac{\Psi_i(v) \Phi_j(\mu)}{\mu - v} \right) \Big|_{\mu=\mu_i, v=v_j},$$

a d

$$Q_1 = \begin{pmatrix} Q^{[11]} & Q^{[12]} & \dots & Q^{[1r]} & \widehat{\chi}_1 \\ Q^{[21]} & Q^{[22]} & \dots & Q^{[2r]} & \widehat{\chi}_2 \\ \vdots & \vdots & \ddots & \vdots & \vdots \\ Q^{[r1]} & Q^{[r2]} & \dots & Q^{[rr]} & \widehat{\chi}_r \\ \widehat{\phi}_1 & \widehat{\phi}_2 & \dots & \widehat{\phi}_r & 0 \end{pmatrix}, \quad Q = \begin{pmatrix} Q^{[11]} & Q^{[12]} & \dots & Q^{[1r]} \\ Q^{[21]} & Q^{[22]} & \dots & Q^{[2r]} \\ \vdots & \vdots & \ddots & \vdots \\ Q^{[r1]} & Q^{[r2]} & \dots & Q^{[rr]} \end{pmatrix},$$

th

$$\widehat{\chi}_i = \left( \chi_i, \frac{\partial}{\partial v} \chi_i, \dots, \frac{1}{(m_i - 1)!} \frac{\partial^{m_i-1}}{\partial v^{m_i-1}} \chi_i \right)^T \Big|_{v=v_i},$$

$$\widehat{\phi}_j = \left( \phi_j, \frac{\partial}{\partial \mu} \phi_j, \dots, \frac{1}{(m_j - 1)!} \frac{\partial^{m_j-1}}{\partial \mu^{m_j-1}} \phi_j \right) \Big|_{\mu=\mu_j},$$

$$Q^{[ij]} = \left( Q_{kl}^{[ij]} \right)_{m_k, m_l},$$

$$Q_{kl}^{[ij]} = -\frac{1}{(k-1)!(l-1)!} \frac{\partial^{k+l-2}}{\partial v^{k-1} \partial \mu^{l-1}} \left( \frac{\Psi_i(v) \sigma_3 \Phi_j(\mu)}{\mu + v} + \frac{\Psi_i(v) \Phi_j(\mu)}{\mu - v} \right) \Big|_{\mu=\mu_i, v=v_j}.$$

AccVd g t V ab V e the V e , t V d f c t t V e e t h a t h e e d d V  
 (28) a d (29) a e t a d f V gDT-I a d gDT-II. F V the gDT-II, t V e d c e  
 t e (4) t V DNLS, e t d  $\mu_i = v_i^*$ .

#### 4. High-order solutions for DNLS

I t e g a b e V e a a t a d f f e e t a e a V a e e V f V the  
 c h e M V t V . T V c V t d t h V e V t V , a b e V f a V a c h e  
 h a e b e e V V e d c d g I S T , D e g d h V d , H V a ' b e a t h e V  
 a d D a b V (B a c d) d h V d , d c .

W h e c a c a D T V t V b e a c V e e t t V t V c V t d N - V t V  
 V t V , t a V b e d e d e d t V M a t h e h g h - V d e V t V , h c h  
 c V e V d t V t e V e M t h e e e d V c V e f c e t t h e I S T t e V V g  
 ( e e [6] a d t h e e f e e c e t h e e ) . W e h V t h e d V t a t h e g D T  
 d e e d a b V e c a b e a e d t V M a a V V t V f V D N L S . I d e e d ,  
 a a t f V t h e h g h - V d e V t V , a d M e N - V t V V t V a V  
 a e a . I S e d V 4.1 , e c V d e t h e V t V t h t h e V B C . I a t c a ,  
 N - a V a V t V , h g h - V d e a V a V t V a d h g h - V d e V t V a e  
 V e d V t . I S e d V 4.2 , e c V t d t h e h g h - V d e V t V t h N V B C  
 h c h c d e t h e h g h - V d e a V a V t V t h N V B C a d h g h - V d e  
 V g e a e V t V .

##### 4.1. Solutions with VBC

A g D T - I t V a c , e a M a t h e e d M V t V , a e  
 a e a e V t V , N - h a e V t V ( e V d c V t V ) a d N - V t V  
 V t V ( e e [17]). A d d t V a , f e t a e t M t h e V t V V t V , e  
 c a c d a V a V t V [17, 27]. I t h e d V , e c V d e t h e N - a V a  
 V t V c t . A e V t e d V t , t h e a V a V t V a e t h e t c a e t t h e  
 V t V V t V . T h e d f f e e t b e h a V M t h e h g h - V d e a V a V t V  
 a d h g h - V d e V t V a e d c a e d .

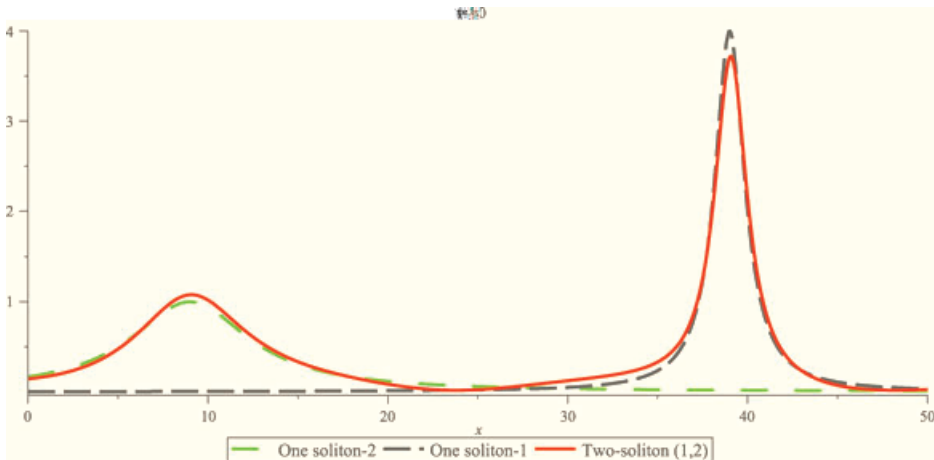
I t h e r t t V c a e , g D T - I b e e d , h e f V t h e c a e 3 , t  
 V e c V e e t t V e t h e g D T - I I c e t h e e d a a a d e e e d t V b e  
 c V g a e d t h e a c h V h e .

##### Case 1: N-rational solutions

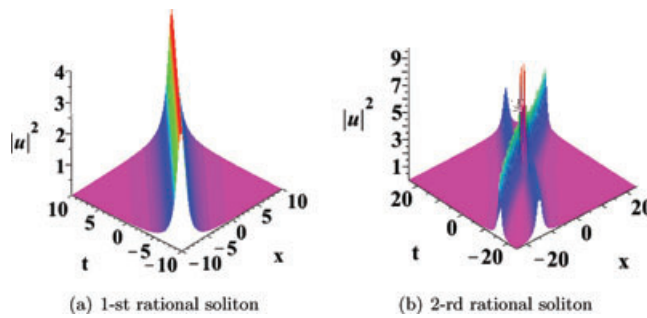
W e c t c V d e t h e a V a V t V a d t h e h g h e V d e a a V g e . F V  
 t h e e e d V t V  $u = 0$ , t h e e c a V t V f V L a a (4) t h t h e e d d V  
 $u = v^*$

$$\begin{pmatrix} \varphi \\ \phi \end{pmatrix} = \begin{pmatrix} e^{-\zeta^{-2}(x+2\zeta^{-2}t+c)} \\ e^{\zeta^{-2}(x+2\zeta^{-2}t+c)} \end{pmatrix}, \quad (38)$$





**Figure 1.** The magnitude of the wave function  $|u|$  versus position  $x$  for the two-soliton solution (1,2) with the parameters  $\zeta_1 = 2$ ,  $\zeta_2 = 4$  and  $c_1 = c_2 = 1$ ; One soliton-1 with the parameters  $\zeta_1 = 4$  and  $c_1 = 1$ ; One soliton-2 with the parameters  $\zeta_1 = 2$  and  $c_1 = 1$ .



**Figure 2.** High-order soliton solution for the VBC: The parameters  $\zeta_1 = 2$  and  $c_1 = 0$ .

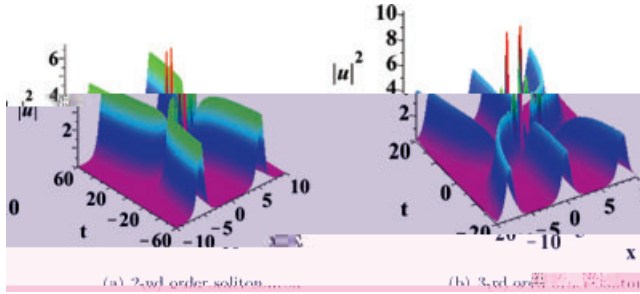
### Case 3: High-order solitons

Take the high-order soliton solution, we take the seed solution  $u = 0$ . The equation (4) can be written as  $u = v^*$  and  $u = 0$

$$\begin{pmatrix} \varphi \\ \phi \end{pmatrix} = \begin{pmatrix} e^{-\mu^{-2}(x+2\mu^{-2}t+c)} \\ e^{\mu^{-2}(x+2\mu^{-2}t+d)} \end{pmatrix},$$

and the equation (5) can be written as  $u = v^*$  and  $u = 0$  and a

$$(\chi, \psi) = (e^{-v^{-2}(x+2v^{-2}t+c^*)}, e^{-v^{-2}(x+2v^{-2}t+d^*)}),$$



**Figure 3.** High-order solitons in the VBC: The parameters  $\mu = 1 + i$  and  $c = d = 0$ .

where  $c, d$  are constants which are determined by the initial data. Then, the NLS equation can be written as

$$u_t[N] = -2 \left( \frac{d}{dt} (M_1) \right)_x, \quad M = (M_{ij})_{N \times N} \quad (41)$$

where

$$M_{ij} = \frac{d^{i+j-2}}{dv^{i-1} d\mu^{j-1}} \left[ \frac{2[v e^{(v^2 - \mu^2)[x + 2(v^2 + \mu^2)t]} + \mu e^{-(v^2 - \mu^2)[x + 2(v^2 + \mu^2)t]}]}{v^2 - \mu^2} \right] \Big|_{v=\mu^*},$$

$$M_1 = \begin{pmatrix} M & Y^T \\ X & 0 \end{pmatrix}, \quad X = \left( \varphi, \frac{d}{d\mu} \varphi, \dots, \frac{d^{N-1}}{d\mu^{N-1}} \varphi \right), \quad Y = \left( \psi, \frac{d}{dv} \psi, \dots, \frac{d^{N-1}}{dv^{N-1}} \psi \right).$$

The evolution of the soliton solution is shown in Figure 3. The high-order solitons in the VBC are shown in Figure 3. The high-order solitons in the VBC are shown in Figure 3. The high-order solitons in the VBC are shown in Figure 3.

#### 4.2. Solution with NVBC

The solution of the NVBC can be obtained by the DT method. A detailed description of the DT method can be found in [26]. The DT method is a powerful tool for solving the NLS equation. The solution of the NLS equation can be obtained by the DT method. The solution of the NLS equation can be obtained by the DT method. The solution of the NLS equation can be obtained by the DT method.

The high-order solitons in the NVBC are shown in Figure 3. The high-order solitons in the NVBC are shown in Figure 3. The high-order solitons in the NVBC are shown in Figure 3. The high-order solitons in the NVBC are shown in Figure 3.

the high-order soliton solution can be obtained by using the ansatz (37) and the boundary conditions (38). The high-order soliton solution can be obtained by using the ansatz (37) and the boundary conditions (38).

### Case 1: High-order rational solitons with NVBC from vacuum

For the high-order soliton solution with NVBC, the boundary conditions are

$$\widehat{Y}_1 = \begin{pmatrix} \widehat{y}_1 \\ \widehat{y}_1^{(1)} \\ \vdots \\ \widehat{y}_1^{(2N)} \end{pmatrix}, \quad \widehat{Z}_1 = \begin{pmatrix} \widehat{z}_1 \\ \widehat{z}_1^{(1)} \\ \vdots \\ \widehat{z}_1^{(2N)} \end{pmatrix},$$

where

$$\widehat{y}_1 = (\varphi, \zeta_1 \phi, \dots, \zeta_1^{2N-1} \phi, \zeta_1^{2N+1} \phi), \quad \widehat{z}_1 = (\varphi, \zeta_1 \phi, \dots, \zeta_1^{2N-1} \phi, \zeta_1^{2N} \varphi),$$

and  $\phi, \psi$  are given by (38). The high-order soliton solution with NVBC can be expressed as

$$u[N] = - \left( \frac{d\widehat{Y}_1}{d\widehat{Z}_1} \right)_x. \quad (42)$$

Taking the limit  $\zeta_1 = a$  (a real constant) and  $c = 0$ , we have the rational soliton solution with NVBC

$$u[1] = \frac{2L_2 L_1^*}{aL_1^2} e^{-\frac{2(a^2 x - 2t)}{a^4}}, \quad (43)$$

where

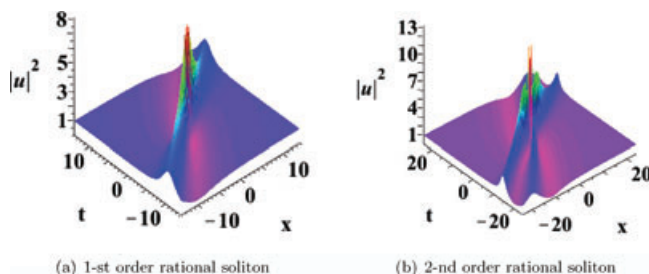
$$L_1 = 16\xi^2 + a^8 + 8a^4(\xi - 4t), \quad L_2 = 16\xi^2 - 3a^8 - 8a^4(\xi + 4t)$$

and  $\xi = a^2 x - 4t$ . The solution  $u[1]$  at the initial time  $t = 0$  is given by, and the phase  $\theta(x, t) = (0, \pm \frac{\sqrt{3}}{16} a^4)$ . The density of the soliton (43) is a function of  $x$  and  $t$  and decays to zero as  $x \rightarrow \pm \infty$ . The rational soliton solution with NVBC is shown in Fig. 4.

### Case 2: High-order rational solitons with NVBC from plane wave solution

For the high-order soliton solution with NVBC, the boundary conditions are

$$u = A e^{-i(2\theta_2)},$$



**Figure 4.** Rational soliton with NVBC: The parameters are  $\zeta_1 = 2$  and  $c_1 = 0$ .

where  $\theta_2 = \frac{1}{2}[ax - (A^2a + a^2)t + c]$ ,  $c \in \mathbb{R}$ . The corresponding data is plotted in Figure 4(b).

$$\Phi = \begin{pmatrix} e^{(\theta_1 + \theta_2) + \phi} & e^{-(\theta_1 - \theta_2) - \phi} \\ e^{-(\theta_1 + \theta_2) + \phi} & e^{(\theta_1 - \theta_2) - \phi} \end{pmatrix}, \quad (44)$$

where

$$\theta_1 = \frac{1}{2}a \coth \left( -\frac{2 + a\zeta^2}{2A\zeta} \right),$$

and

$$\phi = \frac{1}{2}\sqrt{-(2\zeta^{-2} + a)^2 - 4A^2\zeta^{-2}}[x - (a + A^2 - 2\zeta^{-2})t + d].$$

Using the condition (28), we have  $-4A^2\zeta^2 - (2 + a\zeta^2)^2 > 0$ ,  $d \in \mathbb{R}$ ,  $\zeta \in \mathbb{R}$ . The corresponding data is plotted in Figure 4(b). The authors of [26] also studied the case, but the data is not plotted.

If the case  $-4A^2\zeta^2 - (2 + a\zeta^2)^2 = 0$ ,  $\Phi$  given by (44), which does not satisfy the condition (28), so it is not a solution. The corresponding data is plotted in Figure 4(b). The authors of [26] also studied the case, but the data is not plotted.

For the case, we set  $a = c = 0$ ,  $A = 1$  which leads to the case  $-4A^2\zeta^2 - (2 + a\zeta^2)^2 = 0$ . We have the data plotted in Figure 4(b). With the case

$$\Phi_1 = \frac{\Phi|_{\zeta=(1+f)C}}{f^{1/2}}, \quad C = \begin{pmatrix} 1 \\ 1 \end{pmatrix},$$

if  $f = 0$ , which is the case (28), we have

$$\Phi_1 = Y_0 + Y_1f + \cdots + Y_nf^n + \cdots,$$

$$Y_n = \begin{pmatrix} x_n \\ y_n \end{pmatrix} = \lim_{f \rightarrow 0} \frac{1}{n!} \frac{d^n}{df^n} \Phi_1(f).$$

The high-order solutions of the DNLS equation are obtained by the following method:

(i) When  $N = 2l - 1$ , we have

$$u[2l - 1] = -1 - \left[ \frac{d\mathfrak{d}(B)}{d\mathfrak{d}(A)} \right]_x, \quad (45)$$

where

$$A_{i,2j-1} = B_{i,2j-1} = 2^{j-1} \sum_{k=0}^{(i-1,2j-2)} C_{2j-1}^k x_k, \quad (j = 1, 2, \dots, l-1),$$

$$B_{i,2l-1} = 2^l \sum_{k=0}^{i-1} C_{2l}^k x_k,$$

$$A_{i,2j} = B_{i,2j} = 2^j \sum_{k=0}^{(i-1,2j-1)} C_{2j}^k y_k, \quad C_m^n = \frac{m!}{n!(m-n)!}.$$

(ii) When  $N = 2l$ , we have

$$u[2l] = 1 - \left[ \frac{d\mathfrak{d}(D)}{d\mathfrak{d}(C)} \right]_x, \quad (46)$$

where

$$C_{i,2j-1} = D_{i,2j-1} = 2^{j-1} \sum_{k=0}^{(i-1,2j-2)} C_{2j-1}^k x_k,$$

$$C_{i,2j} = 2^j \sum_{k=0}^{(i-1,2j-1)} C_{2j}^k y_k, \quad (j = 1, 2, \dots, l)$$

$$D_{i,2k} = C_{i,2k}, \quad (k = 1, 2, \dots, l-1), \quad D_{i,2l} = 2^{l+1} \sum_{k=0}^{i-1} C_{2l+1}^k y_k.$$

In addition, taking  $d = ef$ , where  $e$  and  $f$  are arbitrary functions, we have

$$x_0 = \sqrt{2}(2x - 6t - 1), \quad y_0 = \sqrt{2}(2x - 6t + 1),$$

$$x_1 = \sqrt{2} \left[ \frac{2}{3}x^3 - 6x^2t + 18xt^2 - 18t^3 - 4x + 20t + 2e + \left( \frac{1}{2} - x^2 + 6xt - 9t^2 \right) \right],$$

$$y_1 = \sqrt{2} \left[ \frac{2}{3}x^3 - 6x^2t + 18xt^2 - 18t^3 - 4x + 20t + 2e + \left( x^2 - \frac{1}{2} - 6xt + 9t^2 \right) \right].$$

Then by using the above formula (45), the first derivative of  $u$  with respect to  $x$  is

$$u[1] = -\frac{(-2x + 6t - 1)(-2x + 6t + 3)}{(-2x + 6t + 1)^2}, \quad (47)$$

which shows that the above formulae are valid for the NVBC. Similarly, the second derivative of  $u$  with respect to  $x$  is

$$u[2] = \frac{L_1^* L_2}{L_1^2}, \quad (48)$$

where

$$L_1 = 8\eta^3 + 18\eta + 48t + 12e + (12\eta^2 + 3),$$

$$L_2 = 8\eta^3 - 30\eta + 48t + 12e + (36\eta^2 - 15),$$

and  $\eta = 3t - x$  and  $e$  is an arbitrary constant. The first derivative of  $u$  with respect to  $x$  is  $u[1] = (-\frac{3}{4}e, -\frac{1}{4}e)$ , and the second derivative of  $u$  with respect to  $x$  is  $u[2] = (\frac{7\tau - 4\tau^3 - 6e}{8}, \frac{15\tau - 4\tau^3 - 6e}{24})$ , where  $\tau = \pm\sqrt{\frac{5}{12}}$ . The first derivative of  $u$  with respect to  $t$  is  $u[1] = (-\frac{3}{4}e, -\frac{1}{4}e)$ , and the second derivative of  $u$  with respect to  $t$  is  $u[2] = (\frac{7\tau - 4\tau^3 - 6e}{8}, \frac{15\tau - 4\tau^3 - 6e}{24})$ . The first derivative of  $u$  with respect to  $x$  is  $u[1] = (-\frac{3}{4}e, -\frac{1}{4}e)$ , and the second derivative of  $u$  with respect to  $x$  is  $u[2] = (\frac{7\tau - 4\tau^3 - 6e}{8}, \frac{15\tau - 4\tau^3 - 6e}{24})$ . The first derivative of  $u$  with respect to  $t$  is  $u[1] = (-\frac{3}{4}e, -\frac{1}{4}e)$ , and the second derivative of  $u$  with respect to  $t$  is  $u[2] = (\frac{7\tau - 4\tau^3 - 6e}{8}, \frac{15\tau - 4\tau^3 - 6e}{24})$ .

### Case 3: High-order rogue wave solutions

The high-order rogue wave solutions of the DNLS equation are derived in [17] and the first derivative of  $u$  with respect to  $x$  is  $u[1] = (-\frac{3}{4}e, -\frac{1}{4}e)$ , and the second derivative of  $u$  with respect to  $x$  is  $u[2] = (\frac{7\tau - 4\tau^3 - 6e}{8}, \frac{15\tau - 4\tau^3 - 6e}{24})$ . The first derivative of  $u$  with respect to  $t$  is  $u[1] = (-\frac{3}{4}e, -\frac{1}{4}e)$ , and the second derivative of  $u$  with respect to  $t$  is  $u[2] = (\frac{7\tau - 4\tau^3 - 6e}{8}, \frac{15\tau - 4\tau^3 - 6e}{24})$ .

The high-order rogue wave solutions of the DNLS equation are derived in [17] and the first derivative of  $u$  with respect to  $x$  is  $u[1] = (-\frac{3}{4}e, -\frac{1}{4}e)$ , and the second derivative of  $u$  with respect to  $x$  is  $u[2] = (\frac{7\tau - 4\tau^3 - 6e}{8}, \frac{15\tau - 4\tau^3 - 6e}{24})$ . The first derivative of  $u$  with respect to  $t$  is  $u[1] = (-\frac{3}{4}e, -\frac{1}{4}e)$ , and the second derivative of  $u$  with respect to  $t$  is  $u[2] = (\frac{7\tau - 4\tau^3 - 6e}{8}, \frac{15\tau - 4\tau^3 - 6e}{24})$ .

$$\Phi = E \begin{pmatrix} \alpha & \alpha^{-1} \\ -\alpha^{-1} & -\alpha \end{pmatrix} \begin{pmatrix} \beta & 0 \\ 0 & \beta^{-1} \end{pmatrix}, \quad E = \begin{pmatrix} e^{-\frac{1}{2}x} & 0 \\ 0 & e^{\frac{1}{2}x} \end{pmatrix}$$

where

$$\mathfrak{d} f=0, \text { a } \mathfrak{d}^{\mathfrak{r}} \mathfrak{d}$$

$$\Phi_1=Y_0+Y_1f+\cdots+Y_nf^n+\cdots,$$

he e

$$Y_n=\left(\begin{smallmatrix}x_n\\y_n\end{smallmatrix}\right)=\quad_{f=0}\frac{1}{n!}\frac{\mathfrak{d}}{\mathfrak{d}f^n}\Phi_1(f).$$

Ect , e ha e

$$x_0=\mathfrak{e}\quad\left[-\frac{1}{2}\,x\right](2x-2\,t-1-\,),$$

$$y_0=\mathfrak{e}\quad\left[\frac{1}{2}\,x\right](2x-2\,t+1+\,),$$

$$x_1=\mathfrak{e}\quad\left[-\frac{1}{2}\,x\right]\left[-\frac{1}{3}x^3+\,x^2t+xt^2-\frac{1}{3}\,t^3+\frac{1+}{2}x^2+(1-\,)xt-\frac{1+}{2}t^2\right.\\ \left.-\frac{1}{2}\,x-\frac{5}{2}x+\frac{13}{2}\,t-\frac{1}{2}t+2e+\frac{1}{2}+2\,g\right],$$

$$y_1=\mathfrak{e}\quad\left[\frac{1}{2}\,x\right]\left[-\frac{1}{3}x^3+\,x^2t+xt^2-\frac{1}{3}\,t^3-\frac{1+}{2}x^2-(1-\,)xt+\frac{1+}{2}t^2\right.\\ \left.-\frac{1}{2}\,x-\frac{5}{2}x+\frac{13}{2}\,t-\frac{1}{2}t+2e-\frac{1}{2}+2\,g\right].$$

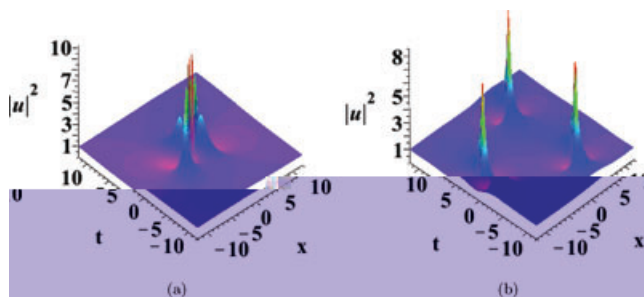
The ef\mathfrak{V}e,the Mh\mathfrak{V}de\mathfrak{V}g e a e ca be e e eted a f\mathfrak{V}\mathfrak{V} g

$$u[N]=\mathfrak{e}\quad[-x]-2\left(\frac{\mathfrak{d}\mathfrak{d}(M_1)}{\mathfrak{d}\mathfrak{d}(M)}\right)_x,\tag{49}$$

he e  $M=(M_{ij})_{N\times N}$ ,

$$M_1=\left(\begin{array}{ccccc}M_{11}&M_{12}&\cdots&M_{1N}&y_0\\M_{21}&M_{22}&\cdots&M_{2N}&y_1\\ \vdots&\vdots&\ddots&\vdots&\vdots\\M_{N1}&M_{N2}&\cdots&M_{NN}&y_{N-1}\\x_0&x_1&\cdots&x_{N-1}&0\end{array}\right),$$

$$M_{ij}=\sum_{k=0,l=0}^{i-1,j-1}-\left(-\frac{1}{2}\right)^{i+j-(k+l+1)}C_{i+j-(k+l+2)}^{i-k-1}\\ \times\big[Y_k^\dagger\sigma_3Y_l(1-\,)^{i-k-1}(1+\,)^{j-l-1}+Y_k^\dagger Y_l(1+\,)^{i-k-1}(1-\,)^{j-l-1}\big].$$



**Figure 5.** SecV d V de Vg e a e ththe a a de (a)  $e = g = 0$ ; (b)  $e = 0$  a d  $g = 100$ .

I at c a , the f t V de a d the ecV d V de Vg e a e V t V a e g e b

$$u[1] = -\frac{[2t^2 + 2x^2 - 3 - 2(x + 3t)][2x^2 + 2t^2 - 2(x - t)]}{[2x^2 + 2t^2 + 1 + 2(x - t)]^2} e^{-x},$$

a d

$$u[2] = \frac{L_1^* L_2}{L_1^2} e^{-x},$$

he e

$$\begin{aligned} L_1 = & 72[e^2 + g^2] + [48x^3 - 144xt^2 + 72x^2 + 144xt - 72t^2 - 144x - 72t \\ & + 36]e + [-144x^2t + 48t^3 + 72x^2 - 144xt - 72t^2 - 72x + 432t \\ & - 36]g + 8x^6 + 24x^4t^2 + 24x^2t^4 + 8t^6 + 24x^5 - 24x^4t + 48x^3t^2 \\ & - 48x^2t^3 + 24xt^4 - 24t^5 - 12x^4 + 48x^3t - 216x^2t^2 + 48xt^3 + 180t^4 \\ & + 48x^3 - 288xt^2 - 336t^3 + 90x^2 - 72xt + 666t^2 + 54x - 198t + 9, \\ L_2 = & 72[e^2 + g^2] + [48x^3 - 144xt^2 - 72x^2 + 432xt + 72t^2 + 144x + 216t \\ & - 180]e + [-144x^2t + 48t^3 + 216x^2 + 144xt - 216t^2 + 216x + 144t \\ & + 36]g + 8x^6 + 24x^4t^2 + 24x^2t^4 + 8t^6 - 24x^5 - 72x^4t - 48x^3t^2 \\ & - 144x^2t^3 - 24xt^4 - 72t^5 - 60x^4 - 144x^3t - 504x^2t^2 - 144xt^3 \\ & - 60t^4 + 48x^3 + 288x^2t + 576xt^2 - 528t^3 - 198x^2 + 504xt - 486t^2 \\ & + 90x + 414t + 45, \end{aligned}$$

e a d f a e a b t a ea be . The e V t V th d ffe e t a a de a e V ed F g. 5.

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